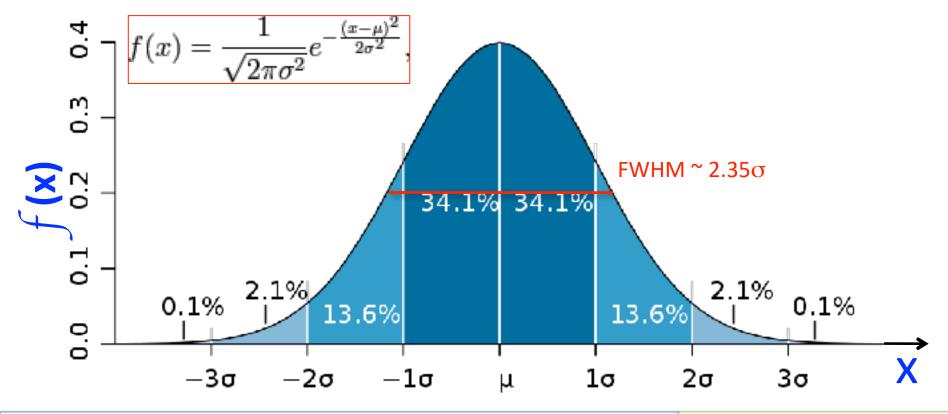
## Lecture 2:

## **Trials and Errors**

- Properties of Normal Distributions
- Trials and Tribulations!
- Regression to the Mean
- Correlations
- Uncertainties & Error Propagation



#### "Two-Sided"

Chance Prob. To Be Within:

 $\pm 1 \sigma$  : 0.6827  $\pm 2 \sigma$  : 0.9545  $\pm 3 \sigma$  : 0.9973  $\pm 4 \sigma$  : 0.99994 Chance Prob. To Be Outside:  $\pm 1 \sigma$  : 0.3173  $\pm 2 \sigma$  : 0.0455  $\pm 3 \sigma$  : 0.0027  $\pm 4 \sigma$  : 6x10<sup>-5</sup>

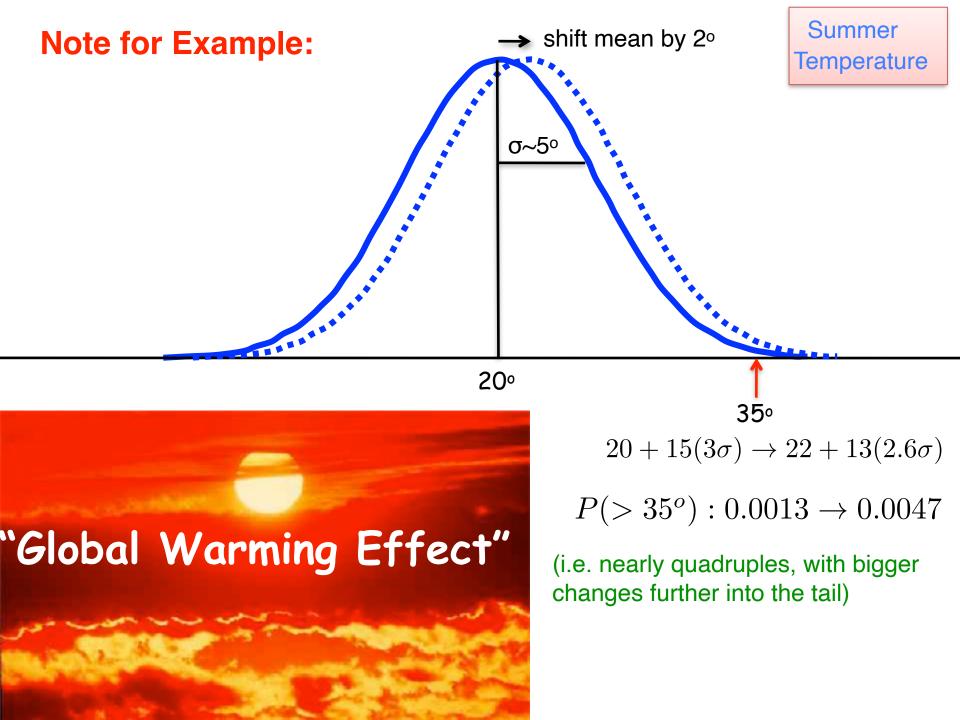
#### e.g. "Do my data points look ok relative to the model?"

#### "One-Sided"

Chance Prob. To Be:

- >1  $\sigma$  : 0.1587
- >3  $\sigma$  : 1.35x10<sup>-3</sup>
- >4 o : 3x10<sup>-5</sup>

*e.g.* "What's the chance of seeing an excess at least this large?"



As a consequence of the Central Limit Theorem, many (but not all!) physical processes often tend towards the Normal Distribution shape.

## However, few achieve this exactly !!

Although calculated probabilities are often couched in terms of an ideal Normal Distribution to give a rough intuition of the scaling

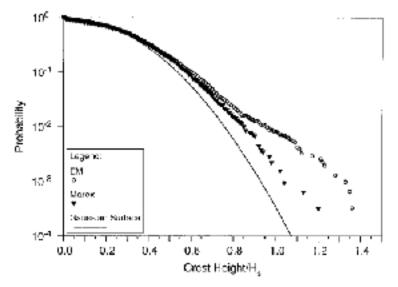
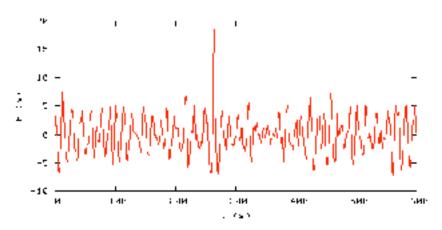


FIG. 4. Probability distribution of normalized crest heights measured at Tern during the storm on 4 Ian 1993. The crest heights are normalized by the significant wave height during each boar of the measurements. Nine hours of measurements with an average significant wave height of about 12 m were combined to produce the observed distribution.

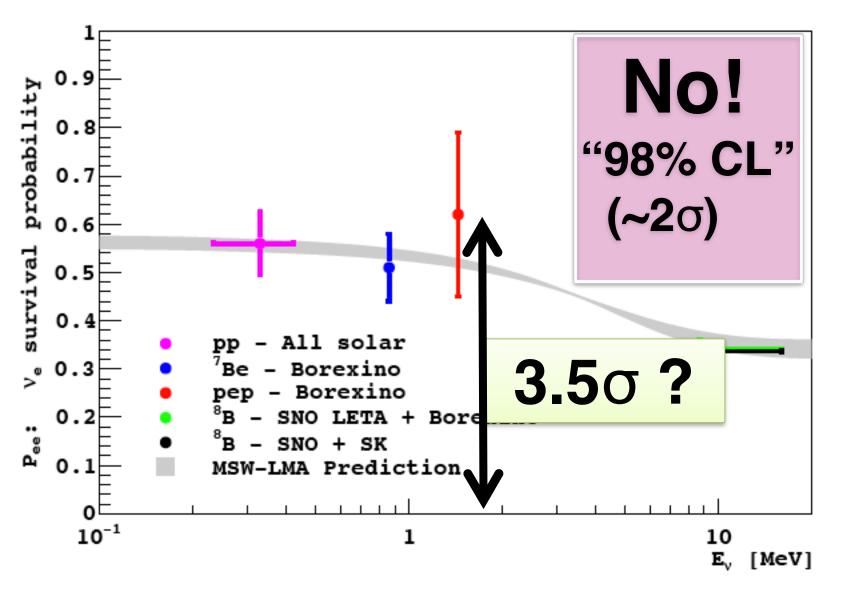


The Draupner wave, a single giant wave measured on New Year's Day 1995, finally confirmed the existence of freak waves, which had previously been considered near-mythical



# Beware the Distribution Tails!!

#### **Borexino "1** $\sigma$ error" on solar pep flux (2011)



## Example: Search for Episodic X-Ray Emission

Over the course of a year, 36000 x-rays are observed to come from a particular astrophysical object. However, on one particular day, 130 events are observed. What is the statistical significance of this observed burst?

$$\langle x \rangle = \frac{36000}{365} = 98.6 \qquad \mu \simeq \langle x \rangle \qquad \sigma = \sqrt{\mu}$$
$$s \cong \frac{(130 - 98.6)}{\sqrt{98.6}} = 3.16\sigma$$

odds of getting at least this many events by a chance fluctuation from the average rate of emission

$$P = 8 \times 10^{-4}$$

Is this sufficient to claim the observation of a burst from this object?



## **Correct question:**

What is the chance of seeing at least one burst with an excess at least as large given the number of independent tests I've done ?

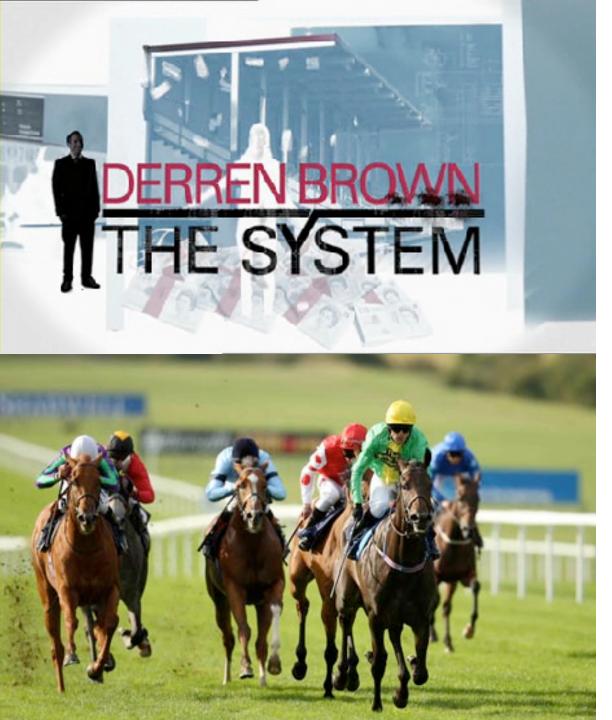
# **Binomial** !!

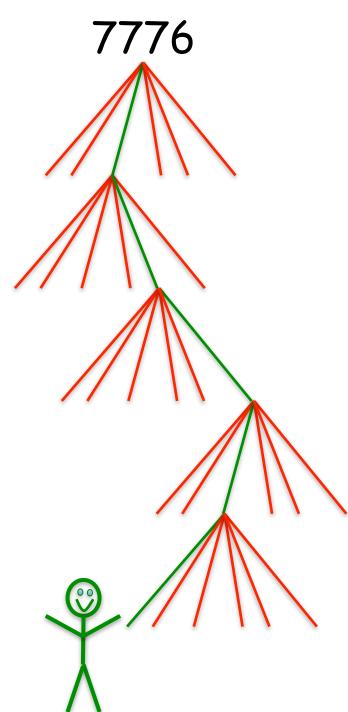
 ${\bf N}$  Bernoulli trials where the chance of each success is  ${\bf P}$ 

$$\sum_{i=1}^{\infty} \binom{N}{i} P^{i} (1-P)^{N-i} = 1 - \binom{N}{0} P^{0} (1-P)^{N-0}$$

$$P_{\text{post-trial}} = 1 - (1 - P)^{N} \quad (\sim NP \quad \text{for } NP <<1)$$
$$P = 8 \ge 10^{-4}, \text{ N} = 365 \quad \longrightarrow \quad P_{\text{post-trial}} = 25\%$$

How many timescales were considered? How many objects examined?





An appreciation of trials factors ("look elsewhere effect") is hugely important... an improper handling of this can lead to incorrect conclusions and opens the door to biased analyses!

### **This is not trivial !** A full accounting for this can be tricky:

- How many hypotheses have you actually tested?
- How many different ways have you tested each hypothesis?
- How many other things would have caught your eye?
- In general, how many ways have you looked at the data?

At the same time, the data needs to be thoroughly checked to look for possible problems and confirm how well it's understood

This is why physicists set the bar high in terms of significance level in order to claim a discovery

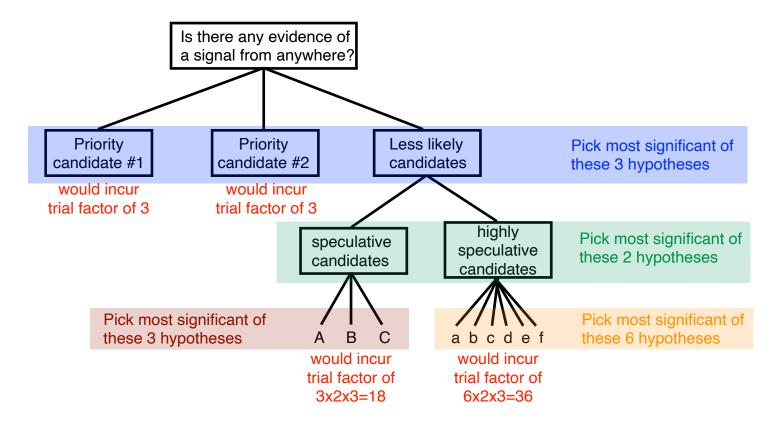
But it's easy to get carried away...



- 1) Trials factors apply to observations that would potentially lead to making a meaningful claim.
- Verification based on applying the same analysis to an independent set of data is a good way to avoid misinterpretation of statistical fluctuations.

# How do you deal with trial factors in the context of an open-ended search when an independent data set may not be available?

It's possible to structure trial factors based on an a priori ranking of hypothesis plausibility\*:



If something is seen from one of these tests that indicates a clear signal, it then belongs to a different population from random fluctuations. So, it can be removed and the search continued with the rest!

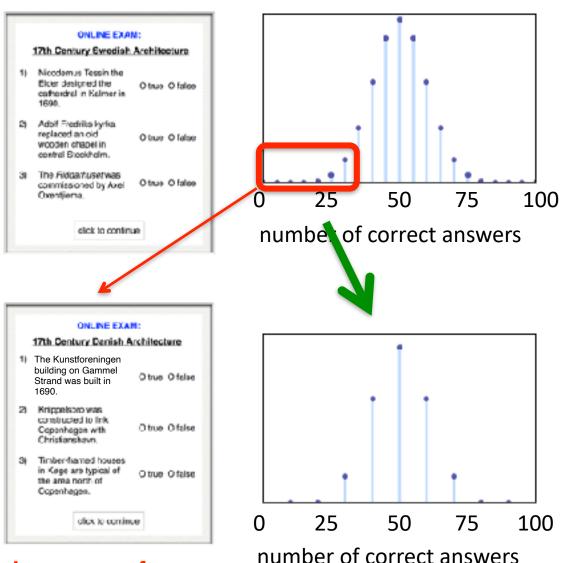
<sup>\*</sup> S. Biller, Astroparticle Physics, Volume 4, Issue 3, p. 285-291

# "Regression to the Mean"

# **Pop Quiz:**

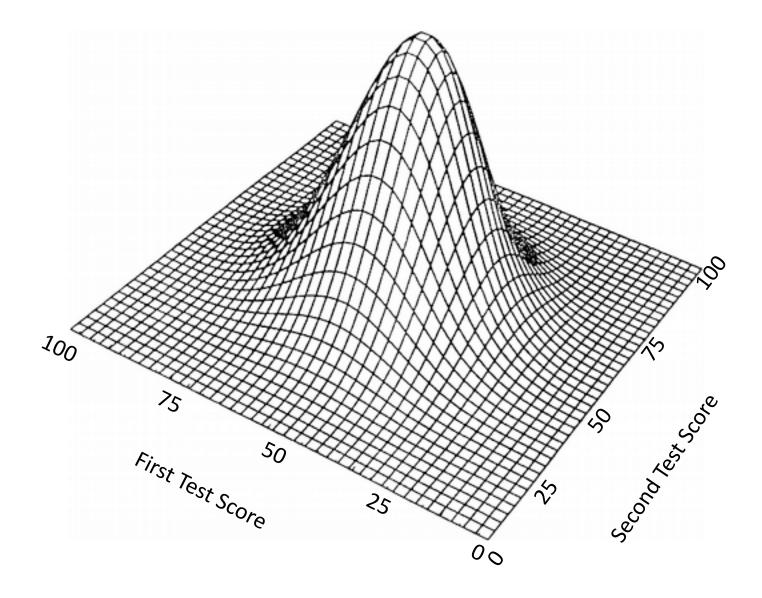
100 true/false questions on 17<sup>th</sup> century Swedish architecture:

100 true/false questions on 17<sup>th</sup> century Danish architecture:

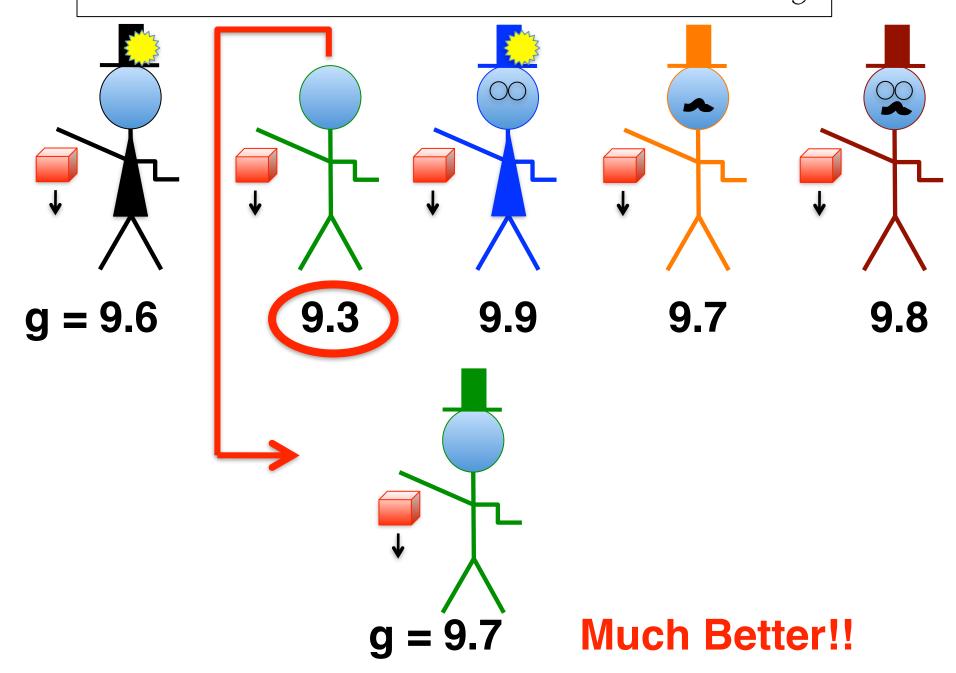


# What an improvement! This particular group of students knows much more about Danish architecture!!

# Bi-variate Distribution with Identical Marginal Distributions (i.e. uncorrelated)

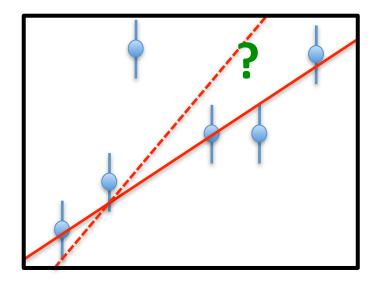


"The Effect of Hats on the Measurement of Gravity"



## So How Do You Handle Outliers?

#### No clear rules!



## **Rules of Thumb:**

- Look for possible systematic biases in the data;
- However, only reject outliers based on clear statistical/scientific criteria;
- Explicitly point out the issue and discuss the details;
- Be aware of any potential bias that could result and review the robustness of your final conclusions.

The total number of known species is ~1.5 million

The number of known species that can fly is ~500,000  $P(flying) > 5x10^{5}/1.5x10^{6} = 0.33$ 

The number of plant species ~400,000 P(plant) =  $4x10^{5}/1.5x10^{6} = 0.27$ 

Thus, probability of finding a flying plant is P(plant) x P(flying) = 0.089

And the expected number of flying plant species is (0.089)(1.5x10<sup>6</sup>) = 133,500

# **Correlations**



Flying	500k 0.35	10k <b>0.007</b>	400 <b>2.8x10-</b> 4		0	Joint
Non-Flying	500k <b>0.35</b>	54 <b>3.8x10</b> -5	6k <b>0.004</b>	400k <b>0.28</b>	10k <b>0.007</b>	PDF
	Insects	Birds	Mammals	Plants	Reptiles	
	500k <b>0.54</b>	54 <b>5.9x10-5</b>	6k <b>0.006</b>	400k <b>0.44</b>	10k <b>0.011</b>	PDF for Non-Flying Species
	Insects	Birds	Mammals	Plants	Reptiles	
	500k <b>0.70</b>	54 <b>0.00704</b>	6k <b>0.00428</b>	400k <b>0.28</b>	10k <b>0.007</b>	"Marginalised" PDF for All Species
	Insects	Birds	Mammals	Plants	Reptiles	

### Just to be clear:

For example, if we have 2 dependent variables, x & y:

$$\int P(x,y)dxdy = 1$$

and

$$\langle f(x,y) \rangle = \int f(x,y) P(x,y) dx dy$$

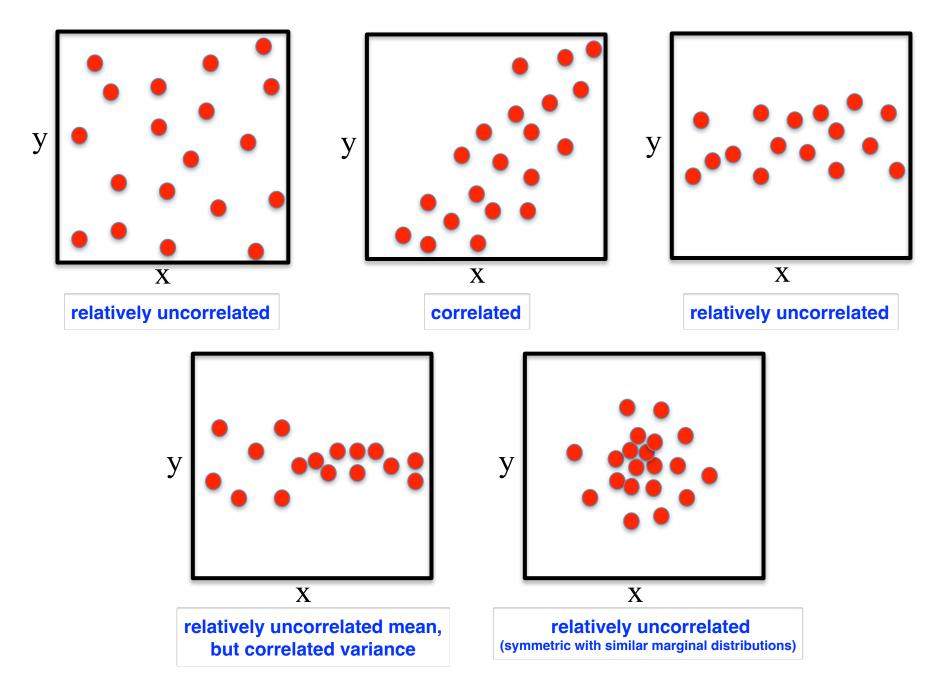
#### How do you tell whether variables are correlated?



# **Correlated or Uncorrelated!**

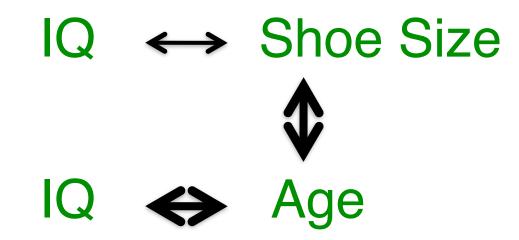


#### **Correlated or Uncorrelated?**





Beware of "hidden" correlations between <u>ANY</u> parameters that distinguish elements of your data set

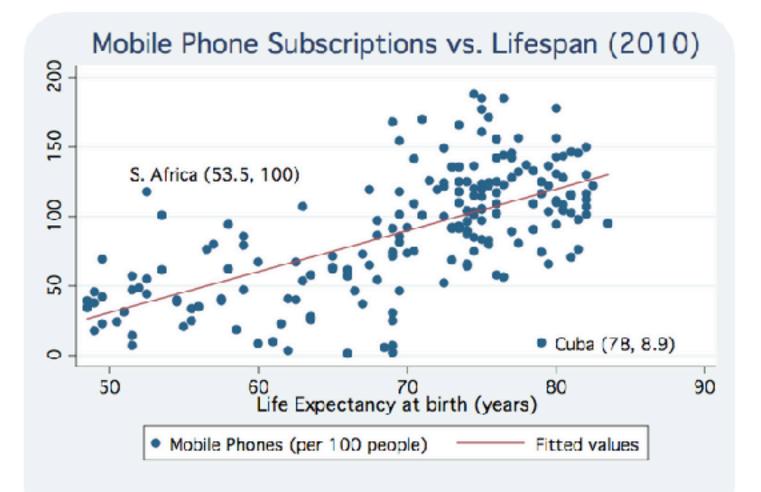






## Beware of jumping to conclusion about cause and effect

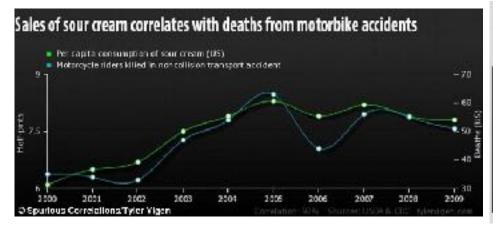




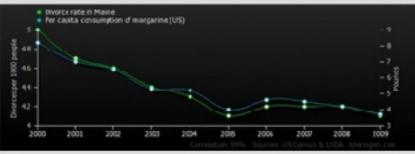


# Beware of spurious correlations

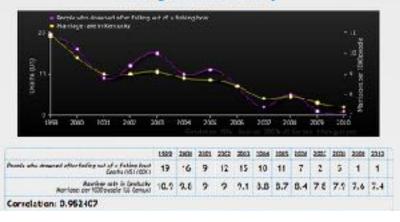




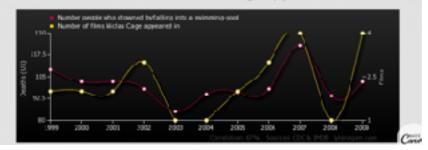
#### Divorce rate in Maine correlates with Per capita consumption of margarine (US)



#### People who drowned after failing out of a fishing boat correctes with Marriage rate in Kentucky

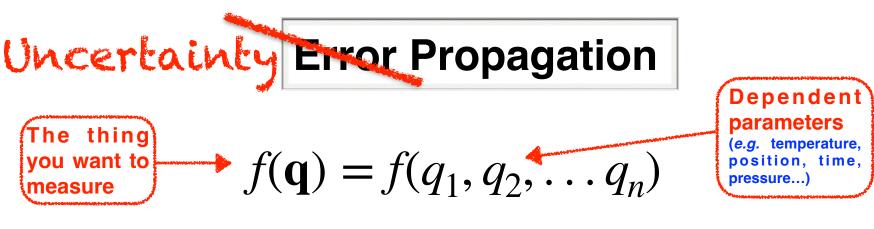


#### Number people who drowned by falling into a swimming-pool computer with Number of films Nicolas Cage appeared in



# Uncertainties





Want to use the distribution f to propagate uncertainties in  $\mathbf{q}$ , but

- 1) We don't necessarily know the full joint distribution of **q** (*i.e.* the probability distribution for all possible sets of values)
- 2) Even if we did, it's cumbersome to deal with!

So, instead, let's approximate things to first order and then estimate the variance of f

$$f(\mathbf{q}) \sim f(\mathbf{\mu}) + \sum_{i=1}^{n} \left[ \frac{\partial f}{\partial q_i} \right]_{\mathbf{q}=\mathbf{\mu}} (q_i - \mu_i)$$

Taylor expansion about the mean values for qwhere  $f(\mathbf{\mu}) = f(\mu_1, \mu_2, \dots, \mu_n)$ 

$$\begin{split} f(\mathbf{q}) \sim f(\mathbf{\mu}) + \sum_{i=1}^{n} \left[ \frac{\partial f}{\partial q_i} \right]_{\mathbf{q}=\mathbf{\mu}} & (q_i - \mu_i) \\ \hline \left\langle f(\mathbf{q}) \right\rangle \sim f(\mathbf{\mu}) + \sum_{i=1}^{n} \left[ \frac{\partial f}{\partial q_i} \right]_{\mathbf{q}=\mathbf{\mu}} & (q_i - \mu_i) = f(\mathbf{\mu}) \\ \hline \left\langle f^2(\mathbf{q}) \right\rangle \sim \left\langle \left[ f(\mathbf{\mu}) + \sum_{i=1}^{n} \left[ \frac{\partial f}{\partial q_i} \right]_{\mathbf{q}=\mathbf{\mu}} (q_i - \mu_i) \right]^2 \right\rangle \\ = f^2(\mathbf{\mu}) + \left\langle \left[ \sum_{i=1}^{n} \left[ \frac{\partial f}{\partial q_i} \right]_{\mathbf{q}=\mathbf{\mu}} (q_i - \mu_i) \right] \left[ \sum_{j=1}^{n} \left[ \frac{\partial f}{\partial q_j} \right]_{\mathbf{q}=\mathbf{\mu}} (q_j - \mu_j) \right] \right\rangle \\ + 2f(\mathbf{\mu}) \sum_{i=1}^{n} \left[ \frac{\partial f}{\partial q_i} \right]_{\mathbf{q}=\mathbf{\mu}} & (q_i - \mu_i) \\ = f^2(\mathbf{\mu}) + \sum_{i,j=1}^{n} \left[ \frac{\partial f}{\partial q_i} \frac{\partial f}{\partial q_j} \right]_{\mathbf{q}=\mathbf{\mu}} & (q_i - \mu_i)(q_j - \mu_j) \end{split}$$

So we get: 
$$\sigma_{f}^{2} = \left\langle f^{2}(\mathbf{q}) \right\rangle - \left\langle f(\mathbf{q}) \right\rangle^{2}$$
$$\simeq \sum_{i,j=1}^{n} \left[ \frac{\partial f}{\partial q_{i}} \frac{\partial f}{\partial q_{j}} \right]_{\mathbf{q}=\mu} \left\langle (q_{i} - \mu_{i})(q_{j} - \mu_{j}) \right\rangle$$
$$= \sum_{i,j=1}^{n} \left[ \frac{\partial f}{\partial q_{i}} \frac{\partial f}{\partial q_{j}} \right]_{\mathbf{q}=\mu} V_{ij}$$

If the q parameters are uncorrelated,

$$V_{ij} = \left\langle (q_i - \mu_i)(q_j - \mu_j) \right\rangle = 0 \quad for \quad i \neq j$$
  
$$V_{ii} = \left\langle (q_i - \mu_i)^2 \right\rangle = \sigma_i^2$$
  
$$\sigma_f^2 \simeq \sum_{i=1}^n \left[ \frac{\partial f}{\partial q_i} \right]_{\mathbf{q}=\mu}^2 \sigma_i^2$$
  
for independent parameter uncertainties

Some simple examples:

$$T_{tot} = t_1 + t_2$$
  
$$\sigma_T^2 \simeq \left[\frac{\partial T}{\partial t_1}\right]^2 \sigma_1^2 + \left[\frac{\partial T}{\partial t_2}\right]^2 \sigma_2^2 = \sigma_1^2 + \sigma_2^2$$

$$s = vt$$
  

$$\sigma_s^2 \simeq \left[\frac{\partial s}{\partial v}\right]^2 \sigma_v^2 + \left[\frac{\partial s}{\partial t}\right]^2 \sigma_t^2 = t^2 \sigma_v^2 + v^2 \sigma_t^2$$
  

$$\left(\frac{\sigma_s}{s}\right)^2 \simeq \left(\frac{\sigma_v}{v}\right)^2 + \left(\frac{\sigma_t}{t}\right)^2$$

For a quadrature addition of uncertainties, uncertainties that are half as big only carry 1/4 of the weight, and uncertainties that are 1/4 as big only carry 1/16 of the weight... **Only the dominant uncertainties matter!** 

#### More General Example: Measurement of Linear Thermal Expansion Coefficient



Measure h by timing the drop of snowballs on one particularly cold day, then compare with  $h_0$ to determine  $\alpha$  In this simple analysis, we're interested in determining the average values of drop time and temperature for the day:

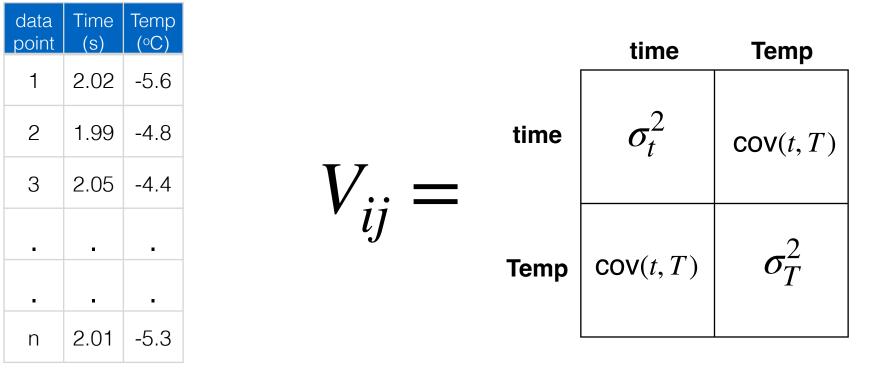
$$T_{\mu} \sim \tilde{t}_{\mu} = \frac{1}{n} \sum_{i=1}^{n} t_i \qquad T_{\mu} \sim \tilde{T}_{\mu} = \frac{1}{n} \sum_{i=1}^{n} T_i$$

Then, from the relation  $h = h_0[1 - \alpha(T - T_0)]$ estimate the expansion coefficient:

$$\tilde{\alpha} = \frac{\frac{\tilde{h_{\mu}}}{h_0} - 1}{\tilde{T_{\mu}} - T_0} = \frac{\frac{g\tilde{t}_{\mu}^2}{2h_0} - 1}{\tilde{T_{\mu}} - T_0}$$

Now we want to find the uncertainty in  $\tilde{\alpha}$  by propagating the uncertainties in  $\tilde{t}_{\mu}$  and  $\tilde{T}_{\mu}$ 

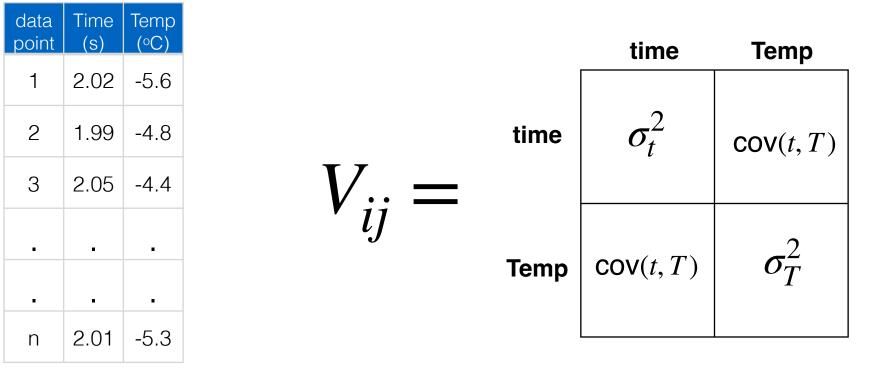
data point	Time (s)	Temp (ºC)
1	2.02	-5.6
2	1.99	-4.8
3	2.05	-4.4
•	•	-
•		-
n	2.01	-5.3



Approach 1: Evaluate Approximately Using the Data

$$\sigma_t^2 = \left\langle (t - t_{\mu})^2 \right\rangle \sim \frac{1}{n - 1} \sum_{i=1}^n (t_i - \tilde{t}_{\mu})^2$$
$$\sigma_T^2 = \left\langle (T - T_{\mu})^2 \right\rangle \sim \frac{1}{n - 1} \sum_{i=1}^n (T_i - \tilde{T}_{\mu})^2$$
$$\operatorname{cov}(t, T) = \left\langle (t - t_{\mu})(T - T_{\mu}) \right\rangle \sim \frac{1}{n - 1} \sum_{i=1}^n (t_i - \tilde{t}_{\mu})(T_i - \tilde{T}_{\mu})$$

Drawback: Requires a large enough data set so that estimates are well determined



Approach 2: Use Calibration Measurements and/or Physical Models

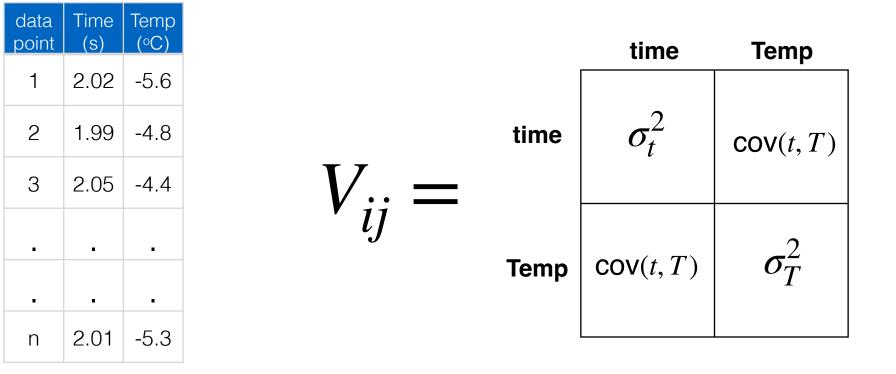
$$\sigma_t^2 = \left\langle (t - t_{\mu})^2 \right\rangle$$
 from calibration of timing accuracy

 $\sigma_T^2 = \left\langle (T - T_\mu)^2 \right\rangle$  from calibration of temperature reading accuracy Temperature variations during the day are sufficiently small that the

correlation with time measurements is very weak, so  $cov(t, T) \sim 0$ 

#### Drawback: Model could be wrong

The temperature variations might influence reaction times and this might have a noticeable systematic impact on the stopwatch measurements



Best Approach: **Do both!** 

Check the consistency or your model and calibrations with the data (If things don't add up, dig around to understand it!)

$$\alpha = \frac{\frac{gt^2}{2h_0} - 1}{T - T_0}$$

$$\left[\frac{\partial \alpha}{\partial t}\right]_{t_\mu, T_\mu} = \frac{\frac{gt_\mu}{h_0} - 1}{T_\mu - T_0}$$

$$\left[\frac{\partial \alpha}{\partial T}\right]_{t_\mu, T_\mu} = -\frac{\frac{gt_\mu^2}{2h_0} - 1}{(T_\mu - T_0)^2}$$

$$\sigma_{\alpha}^{2} = \left( \begin{bmatrix} \frac{\partial \alpha}{\partial t} \end{bmatrix}_{t_{\mu}, T_{\mu}}, \begin{bmatrix} \frac{\partial \alpha}{\partial T} \end{bmatrix}_{t_{\mu}, T_{\mu}} \right) \begin{pmatrix} \sigma_{t}^{2} & \operatorname{cov}(t, T) \\ \operatorname{cov}(t, T) & \sigma_{T}^{2} \end{pmatrix} \begin{pmatrix} \begin{bmatrix} \frac{\partial \alpha}{\partial t} \end{bmatrix}_{t_{\mu}, T_{\mu}} \\ \begin{bmatrix} \frac{\partial \alpha}{\partial T} \end{bmatrix}_{t_{\mu}, T_{\mu}} \end{pmatrix}$$

$$\sigma_{\alpha}^{2} = \left[\frac{\partial\alpha}{\partial t}\right]_{\mu}^{2} \sigma_{t}^{2} + \left[\frac{\partial\alpha}{\partial T}\right]_{\mu}^{2} \sigma_{T}^{2} + 2\left[\frac{\partial\alpha}{\partial t}\frac{\partial\alpha}{\partial T}\right]_{\mu} \operatorname{cov}(t,T)$$

#### The Statistical Calculation That You Should Have Done at the Start!

Typical linear expansion coefficients for building materials  $\sim$ 5x10<sup>-6</sup> per °C Take (T-T<sub>o</sub>) ~ 20°C

 $h_0-h \sim (20m)(5x10^{-6})(20^{\circ}C) = 0.002m$ 

velocity at impact =  $\sqrt{2gh_0} = \sqrt{2(9.8m/s^2)(20m)} \sim 20m/s$ 

So, timing must be known to an accuracy of (0.002/20) = 0.0001s

Accuracy of any one timing timing measurement ~ 0.1s

But we improve by averaging lots of measurements according to  $\sigma_m = \frac{\sigma}{\sqrt{n}}$ 

How many measurements do we need?

$$n = \frac{\sigma^2}{\sigma_m^2} \sim \left(\frac{0.1}{0.0001}\right)^2 = 10^6 \qquad \text{(ignoring system)}$$

(ignoring systematic uncertainties!)

## **Statistical Uncertainties**

**Fundamental, calculable, random variations** due to an inherent limited sampling of the underlying distribution (i.e. counting statistics).

## **Systematic Uncertainties**

**Incidental, estimated (bounded), systematic biases** incurred as a result of limited measurement precision (also always present). There is no universally applicable method for estimating/bounding\* systematic uncertainties. A typical approach often relies on independent cross-checks, accounting for possible statistical limitations of calibration procedures, knowledge about the experimental design and general consistency arguments.

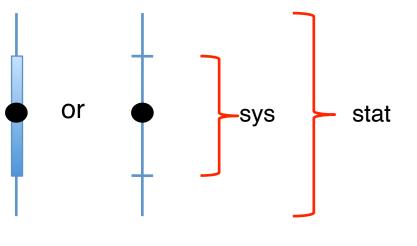
\* Systematic errors that are "determined" become corrections!

Because of their very different nature, there is no standard, mathematically rigorous way to combine the 2 types of uncertainties. The convention is thus to quote results in the form:

## **Result ± Uncertainty (stat) ± Uncertainty (sys)**

And error bars such as:

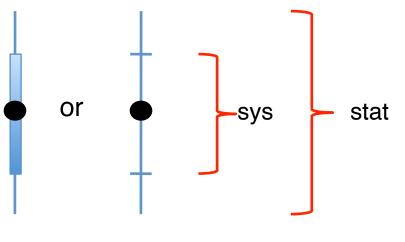
(Much more on error bars later!!)



# How do you then make use of such data points to fit a model?

It is often generally assumed that systematic uncertainties can be treated in a similar way to statistical uncertainties, with careful attention to correlations.

Ideally, the best way to treat systematic uncertainties are as free parameters in the model fit, constrained by the separately determined bounds on their values.



#### (Much more on error bars later!!)