

## Part B Electromagnetism: Problem Sheet 3

[The material on this problem sheet covers sections 3 and 4 of the lecture notes.]

1. A point particle of mass  $m$  and charge  $q$  moves with velocity  $\dot{\mathbf{r}} = \mathbf{v}$  in a constant magnetic field  $\mathbf{B} = B \mathbf{e}_3$ , where  $\mathbf{e}_3$  is the unit vector parallel to the  $z$ -axis and  $B$  is a constant.
  - (a) Write down Newton's second law for the particle, assuming the force on the particle is  $q \mathbf{v} \wedge \mathbf{B}$ , and show that  $\mathbf{v} \cdot \mathbf{e}_3$  and  $v \equiv |\mathbf{v}|$  are constants of the motion.
  - (b) Show that if  $\mathbf{v} \cdot \mathbf{e}_3 = 0$  then the particle follows a circular path in a plane of constant  $z$ , with speed  $v$  and radius  $a$  related by  $mv = aq |\mathbf{B}|$ . What is the angular frequency of this circular motion?
  - (c) What is the path if  $\mathbf{v} \cdot \mathbf{e}_3 \neq 0$ ?
2. Consider a circular loop of wire of radius  $a$  carrying a steady current  $I$ . If the loop is placed in the  $(x, y)$ -plane  $\{z = 0\}$ , centred at the origin, show using the integral form of the Biot-Savart law that the magnetic field at the point  $(0, 0, z)$  on the axis of symmetry is  $\mathbf{B} = (0, 0, B(z))$ , where

$$B(z) = \frac{\mu_0 I a^2}{2(a^2 + z^2)^{3/2}} .$$

3. Consider an infinitely long cylindrical conductor of radius  $a$ , its axis of symmetry being the  $z$ -axis, carrying a uniform current in the  $z$ -direction of constant density  $J$ . Using Ampère's law, and assuming that the magnetic field is of the form  $\mathbf{B} = B(s) \mathbf{e}_\varphi$  in cylindrical polar coordinates  $(s, \varphi, z)$ , show that

$$B(s) = \begin{cases} \frac{\mu_0 J s}{2}, & s \leq a \\ \frac{\mu_0 J a^2}{2s}, & s \geq a . \end{cases}$$

Hence show that the magnetic field outside the conductor is the same as that (derived in lectures) for an infinitely long wire carrying the same current  $I$ .

4. (a) In Cartesian coordinates in *two dimensions* we define  $V_x \equiv -\partial G / \partial y$ ,  $V_y \equiv \partial G / \partial x$ , where  $G = G(x, y)$ . Show using *Green's theorem* that

$$\int_C V_x dx + V_y dy = \int_D \left( \frac{\partial^2 G}{\partial x^2} + \frac{\partial^2 G}{\partial y^2} \right) dx dy ,$$

where  $C$  is a positively oriented simple closed curve bounding a region  $D \subset \mathbb{R}^2$ .

On  $\mathbb{R}^2 \setminus \{(0, 0)\}$  we next define

$$G(x, y) \equiv \frac{1}{2\pi} \log \sqrt{x^2 + y^2} .$$

Show that  $G$  satisfies the Laplace equation on  $\mathbb{R}^2 \setminus \{(0, 0)\}$ , but also

$$\int_C V_x dx + V_y dy = 1 ,$$

where  $C$  is a circle of radius  $a > 0$  centred on the origin. Deduce that

$$\left( \frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2} \right) G = \delta(x) \delta(y) .$$

- (b) Now in *three dimensions* define the magnetostatic potential  $\mathbf{A} \equiv (0, 0, -kG(x, y))$ , with  $G(x, y)$  given in part (a). Show that  $\mathbf{A}$  satisfies the Lorenz gauge condition, together with the magnetostatic Maxwell equation

$$\nabla^2 \mathbf{A} = -\mu_0 \mathbf{J} ,$$

where you should determine the current  $\mathbf{J}$ . Have you seen this problem before? Show that the magnetic field  $\mathbf{B} \equiv \nabla \wedge \mathbf{A} = k(V_x, V_y, 0)$ , and hence show that

$$\int_C \mathbf{B} \cdot d\mathbf{r} = k ,$$

for any closed curve  $C$  which winds once around the  $z$ -axis anticlockwise.

[*Hint*: First take  $C$  to be a circle in the plane  $\{z = 0\}$  centred at the origin, then use Stokes' theorem to obtain the general result.]

5. Assume  $r \neq 0$  throughout this question. Define the vector field  $\mathbf{A}$  by

$$\mathbf{A} \equiv \frac{1}{r^3} \mathbf{k} \wedge \mathbf{r} ,$$

where  $\mathbf{k}$  is a constant vector, and  $r = |\mathbf{r}|$ , as usual.

- (a) Show that  $\mathbf{A}$  can also be written as  $\mathbf{A} = \nabla\phi \wedge \mathbf{k}$  with  $\phi = 1/r$ , and hence show that  $\nabla \cdot \mathbf{A} = 0$ .  
 (b) Defining the magnetic field  $\mathbf{B} \equiv \nabla \wedge \mathbf{A}$ , show that  $\mathbf{B}$  can be written as

$$\mathbf{B} = \nabla [\mathbf{k} \cdot \nabla(1/r)] .$$

Deduce that  $\nabla \cdot \mathbf{B} = 0$ , and that  $\nabla \wedge \mathbf{B} = \mathbf{0}$ .

- (c) Show that

$$\mathbf{B} = \frac{3\mathbf{k} \cdot \mathbf{r}}{r^5} \mathbf{r} - \frac{\mathbf{k}}{r^3} ,$$

which thus describes the magnetic field of a *magnetic dipole*, with  $\mathbf{k} = (\mu_0/4\pi) \mathbf{m}$  where  $\mathbf{m}$  is the magnetic dipole moment. Hence show that

$$\lim_{r \rightarrow \infty} \int_{\Sigma_r} \mathbf{B} \cdot d\mathbf{S} = 0 ,$$

where  $\Sigma_r$  is a sphere of radius  $r$  centred at the origin. Using these results, explain why the magnetic flux through *any* closed surface  $\Sigma$  is zero. (The magnetic dipole field hence has zero magnetic charge, even at the origin where it is singular.)

6. (*Optional*) Show that an electric dipole at the origin, with dipole moment  $\mathbf{p}_0$ , exerts a force  $\mathbf{F}$  on another electric dipole at position  $\mathbf{r}$ , with dipole moment  $\mathbf{p}$ , given by

$$\mathbf{F} = \frac{3}{4\pi\epsilon_0 r^4} [(\mathbf{p} \cdot \hat{\mathbf{r}}) \mathbf{p}_0 + (\mathbf{p}_0 \cdot \hat{\mathbf{r}}) \mathbf{p} + (\mathbf{p}_0 \cdot \mathbf{p}) \hat{\mathbf{r}} - 5(\mathbf{p}_0 \cdot \hat{\mathbf{r}})(\mathbf{p} \cdot \hat{\mathbf{r}}) \hat{\mathbf{r}}] .$$

Here as usual  $\hat{\mathbf{r}} \equiv \mathbf{r}/r$ , with  $r \equiv |\mathbf{r}|$ .

(The same formula holds for a pair of magnetic dipoles *e.g.* a pair of iron filings, replacing electric dipole moments by magnetic dipole moments, and  $\epsilon_0$  by  $1/\mu_0$ .)

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